

Commutative nonassociative algebras, representations of finite groups and minimal cones

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Introduction

The most straightforward approach to an algebra (like \mathbb{C} or \mathbb{H}) is to describe it by a multiplication table. Sometimes one defines an algebra via an identity (like Lie and Jordan algebras). But, if you go beyond small dimensions, this approach is not very satisfactory.

Let A be a commutative nonassociative algebra over \mathbb{F} decomposed into a direct sum

$$A = A_1 \oplus A_2 \oplus \dots \oplus A_m.$$

The most interesting case is when A_i are invariant subspaces of the adjoint operator

$$\text{ad}_c : x \rightarrow cx, \quad c \text{ being an idempotent in } A.$$

A *fusion law* or *fusion rule* is a map $\star : (i, j) \rightarrow i \star j$ such that

$$A_i A_j \subset \bigoplus_{k \in i \star j} A_k.$$

A fusion law is 'nice' if both m and the cardinality of $i \star j$ are reasonably small. There are only few known examples with nice (i.e. graded) fusion laws. All of them have important features and connections to different areas of mathematics.

A key example: Jordan algebra of symmetric matrices

A commutative algebra A is **Jordan** if $x^2(xy) = x(x^2y)$ for all $x, y \in A$.

Let A be the 3-dimensional Jordan algebra of real symmetric 2×2 -matrices with multiplication

$$x \bullet y = \frac{1}{2}(xy + yx)$$

Note that $x \bullet x = xx = x^2$. Set $c = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$, $e_2 = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$, $e_3 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$. Then

\bullet	e_1	e_2	e_3
e_1	e_1	0	$\frac{1}{2}e_3$
e_2	0	e_2	$\frac{1}{2}e_3$
e_3	$\frac{1}{2}e_3$	$\frac{1}{2}e_3$	$e_1 + e_2$

$\Rightarrow \{e_1, e_2, e_3\}$ is an eigenbasis of ad_c and

$$A = A_1 \oplus A_0 \oplus A_{\frac{1}{2}} \text{ with}$$

$A_1 = \langle e_1 \rangle$, $A_0 = \langle e_2 \rangle$ and $A_{\frac{1}{2}} = \langle e_3 \rangle$.

This example illustrates an important general fact: A satisfies the *Jordan-type* fusion law:

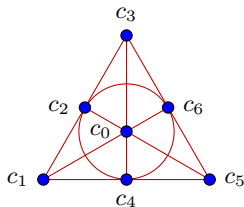
\star	1	0	$\frac{1}{2}$
1	1	\emptyset	$\frac{1}{2}$
0	\emptyset	0	$\frac{1}{2}$
$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1, 0

Finite simple groups

Just as prime numbers, finite simple groups are the building blocks of finite groups.

The Classification (1981): any finite simple group occurs either in one of 18 regular infinite series like cyclic or Lie type groups, or it is one of the 26 *sporadic groups*.

- **Example.** The linear group $L_2(7) \cong L_3(2)$ is the second smallest non-abelian simple group (of order 168). It is the group of automorphisms of the Fano plane.

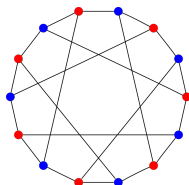


The Fano plane

1	123
2	145
3	160
4	246
5	250
6	340
0	356

points

lines



Incidence graph
(Heawood graph)

- $L_2(7)$ is a $2A$ -generated subgroup of the Monster, the largest sporadic group.

A mystery of the Monster sporadic group

- The Monster Group \mathbb{M} is the largest of the *sporadic* simple groups: its order is

$$2^{46} \cdot 3^{20} \cdot 5^9 \cdot 7^6 \cdot 11^2 \cdot 13^3 \cdot 17 \cdot 19 \cdot 23 \cdot 29 \cdot 31 \cdot 41 \cdot 47 \cdot 59 \cdot 71 \cdot \approx 8 \cdot 10^{53}$$

- **196883** is the dimension of the first nontrivial representation of the Monster.
- In 1978 McKay noticed that the Fourier coefficients of the modular function

$$\begin{aligned} j(\tau) &= \frac{(\theta_2^3(\tau) + \theta_3^3(\tau) + \theta_4^3(\tau))^3}{8\eta(\tau)^{24}} = \sum_{n=-1}^{\infty} c_n q^n \\ &= q^{-1} + 744 + (196883 + 1)q + \dots, \quad q = e^{2\pi i\tau} \end{aligned}$$

are linear combinations of the character degrees of the Monster! McKay and Thompson suggested that a natural graded module for \mathbb{M} might exist.

- Conway and Norton formulate 'MONSTROUS MOONSHINE' (1979).
- In 1988, Frenkel, Lepowsky & Meurman *explicitly* constructed a graded **vertex operator algebra** representation

$$V^\sharp = \bigoplus_{n \geq 0} V_n \quad \text{of } \mathbb{M} \text{ with } V_2 = V_{\mathbb{M}}.$$

- In 1992 Borcherds finally settled the Moonshine conjecture (Fields medal 1998)

Griess algebra

- The Monster was constructed in 1982 by Griess as a group of automorphisms of a **non-associative commutative unital algebra** V_M of dimension $1 + 196883$ over \mathbb{R} , now called the Griess algebra.
- The Griess algebra admits a natural invariant associating bilinear form, i.e.

$$\langle xy, z \rangle = \langle x, yz \rangle, \quad \forall x, y, z \in V_M$$

- The Monster is generated by its involutions (the so-called (2A)-involutions).
- There is a natural bijection between the (2A)-involutions and a subset of idempotents (axes) in V_M
- Any axis induces an orthogonal decomposition $A = A_1 \oplus A_0 \oplus A_{\frac{1}{4}} \oplus A_{\frac{1}{32}}$ with the **Ising fusion laws**:

*	1	0	$\frac{1}{4}$	$\frac{1}{32}$
1	1	0	$\frac{1}{4}$	$\frac{1}{32}$
0	0	0	$\frac{1}{4}$	$\frac{1}{32}$
$\frac{1}{4}$	$\frac{1}{4}$	$\frac{1}{4}$	1, 0	$\frac{1}{32}$
$\frac{1}{32}$	$\frac{1}{32}$	$\frac{1}{32}$	$\frac{1}{32}$	1, 0, $\frac{1}{4}$

*	+	+	+	-
+	+	+	+	-
+	+	+	+	-
+	+	+	+	-
-	-	-	-	+

Griess algebra as a Majorana algebra

*	1	0	$\frac{1}{4}$	$\frac{1}{32}$
1	1	0	$\frac{1}{4}$	$\frac{1}{32}$
0	0	0	$\frac{1}{4}$	$\frac{1}{32}$
$\frac{1}{4}$	$\frac{1}{4}$	$\frac{1}{4}$	1, 0	$\frac{1}{32}$
$\frac{1}{32}$	$\frac{1}{32}$	$\frac{1}{32}$	$\frac{1}{32}$	1, 0, $\frac{1}{4}$

*	+	+	+	-
+	+	+	+	-
+	+	+	+	-
+	+	+	+	-
-	-	-	-	+

- $A_+ := A_1 \oplus A_0 \oplus A_{\frac{1}{4}}$ is a subalgebra
- There is a natural grading $A = A_+ \oplus A_-$ with $A_- := A_{\frac{1}{32}}$
- A key element of the construction: the Majorana involution

$$x \rightarrow x^\tau := x_1 + x_0 + x_{\frac{1}{4}} - x_{\frac{1}{32}}$$

- Furthermore, if $x_{\frac{1}{32}} = 0$ then $x \rightarrow x^\sigma := x_1 + x_0 - x_{\frac{1}{4}}$ is also an algebra involution.

The above algebra structure determines the Monster as an isomorphism group.

G-NA-VOA project

- In his talks on G-NA-VOA project, Robert Griess pointed out that a good theory of commutative nonassociative algebras generalizing the Monster algebra is highly relevant to developing a unified framework of finite simple groups, especially the sporadic groups.
- A starting point is the commutative nonassociative algebra $V_2 = V_{\mathbb{M}}$ in the VOA V^{\sharp} .
- Via foundational work of M.Miyamoto (1995) and later by A.A.Ivanov (2007), the Griess algebra V_2 has been understood as a commutative non-associative algebra generated by a finite set of *primitive* idempotents (i.e. the dimension of the 1-eigenspace is one) satisfying certain axioms, the so-called *Majorana algebras*.
- An algebra is *Majorana* if it possesses a positive definite associating bilinear form $\langle x, y \rangle$ and it is generated (as an algebra) by a finite set of idempotents with the $(1, 0, \frac{1}{4}, \frac{1}{32})$ -Peirce decomposition and the Ising fusion laws. Some further (somewhat subtle) axioms like the Norton inequality are also required.
- (Ivanov, *The Monster Group and Majorana Involutions*, CTM, 2009) proves that the Griess algebra of the Monster is a Majorana algebra and that the action of the Monster on its algebra realizes a Majorana representation of the Monster.

Axial algebras

- A further, radical, step was taken by J. Hall, S. Shpectorov & F. Rehren (2015), they put the previous work into a natural context of *axial algebras*.
- An *axial algebra* over the field \mathbb{F} is a commutative algebra generated by idempotents whose adjoint action has the same multiplicity-free minimal polynomial.

Axial algebras include Jordan algebras, the Matsuo algebras for groups of 3-transpositions, as well as the 196884-dimensional Griess algebra.

- Recall that V_M is generated by $(2A)$ -idempotents satisfying the Ising fusion laws. J. Hall, S. Shpectorov and F. Rehren, (2015) completely characterized axial algebras satisfying the Jordan type fusion laws for $\eta \neq \frac{1}{2}$

\star	1	0	η
1	1	\emptyset	η
0	\emptyset	0	η
η	η	η	1, 0

The case $\eta = \frac{1}{2}$ is singular (contains also Jordan algebras) and is not settled yet.

How incident is a particular Peirce spectrum?

We have seen that

- for Jordan algebras $\mathbf{spec}(c) \subset \{1, 0, \frac{1}{2}\}$
- while for the Griess algebra $\mathbf{spec}(c) \subset \{1, 0, \frac{1}{4}, \frac{1}{32}\}$

(Krasnov Ya., V.T., 2018) proved that a *generic* algebra does not contain $\frac{1}{2}$ in its spectrum. On the other hand, we have

Theorem (V.T., 2018, submitted)

Let A be a finite dimensional commutative nonassociative algebra over a field of characteristic $\neq 2, 3$ and let A satisfy a nontrivial identity $P(z) = 0$. Then

- (A) $\frac{1}{2} \in \mathbf{spec}(c)$ for any idempotent $c \neq 0$;
- (B) if c is semi-simple and λ is single root of the Peirce polynomial $\varrho(P, t)$ then

$$A_\lambda A_{\frac{1}{2}} \subset A_\lambda^\perp := \bigoplus_{\substack{\nu \in \mathbf{spec}(c) \\ \nu \neq \lambda}} A_\nu.$$

In particular, $A_{\frac{1}{2}} A_{\frac{1}{2}} \subset A_{\frac{1}{2}}^\perp$.

A toy example: the Harada algebra

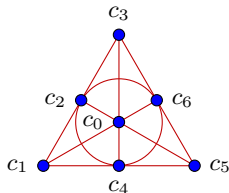
Motivated by Griess' construction of the Monster, the same year Koichiro Harada defined a 6-dimensional algebra H whose automorphism group is $L_2(7)$ as follows:

c_1, \dots, c_6 generate H as a vector space with multiplication

$$c_i c_j = \begin{cases} c_i & \text{if } i = j \\ c_{i \wedge j}, & \text{if } i \neq j, \end{cases}$$

and an additional identity

$$c_0 + c_1 + \dots + c_6 = 0.$$



For example, $c_1 c_2 = c_{1 \wedge 2} = c_3$.

- Each line is a subalgebra!
- Any automorphism $g \in \text{Aut}(H)$ preserves idempotents: $x^2 = x \Rightarrow g(x)^2 = g(x)$.
- Hence, $\text{Aut}(H)$ acts on Φ by permutations.
- Since $\{c_i\}$ generates H , the group $\text{Aut}(H)$ acts faithfully on Φ .
- Since each line is a subalgebra, $\text{Aut}(H)$ act naturally on lines.

This proves that

$$\text{Aut}(H) = \text{Aut}(\Phi) = L_2(7).$$

The Harada algebra

- Φ has totally

$$\#\text{Idem}(H) = 2^{\dim H} - 1 = 63 = 7 + 7 + 21 + 28 \quad \text{nonzero idempotents.}$$

- There are exactly five orbits of the $\text{Aut}(H)$ -action on $\text{Idem}(H)$:

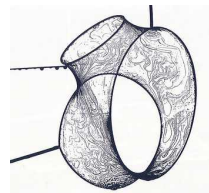
- ▶ 7 points: c_i
- ▶ 7 lines: $\frac{1}{3}(c_{i \wedge j} + c_i + c_j)$,
- ▶ 21 flags: $\frac{1}{3}(2c_{i \wedge j} - c_i - c_j)$;
- ▶ 28 anti-flags: $-3(c_{i \wedge j} + c_i + c_j) - 5c_m$.

Theorem (Ya. Krasnov, T., 2018, in preparation) Idempotents in each orbit have the same spectrum and the only graded (Ising type) fusion laws occur for flags:

*	1	0	-1	$\frac{4}{3}$	$-\frac{2}{3}$
1	1	•	-1	$\frac{4}{3}$	$-\frac{2}{3}$
0	•	0	•	$\frac{4}{3}$	$-\frac{2}{3}$
-1	-1	•	1	•	$-\frac{2}{3}$
$\frac{4}{3}$	$\frac{4}{3}$	$\frac{4}{3}$	•	1, 0	$-\frac{2}{3}$
$-\frac{2}{3}$	$-\frac{2}{3}$	$-\frac{2}{3}$	$-\frac{2}{3}$	$-\frac{2}{3}$	1, 0, $\frac{1}{4}$

Algebras of cubic minimal cones

A **minimal surface** minimizes locally the area functional. Geometrically, this means that the **mean curvature** $= 0$. Analytically, a solution of certain nonlinear PDE.



A **minimal cone** is a typical singularity of a minimal surface. The most of known minimal cones are algebraic, i.e. zero level sets of a homogeneous polynomial $u \in \mathbb{R}[x_1, \dots, x_n]$.

Degree 2 minimal cones are well-known. An example is the Clifford-Simon cone given by

$$u(x) := (x_1^2 + x_2^2 + x_3^2 + x_4^2) - (x_5^2 + x_6^2 + x_7^2 + x_8^2)$$

which is also known as the norm for split octonions. This cone played a crucial role in the solution of the famous Bernstein problem and counterexamples in higher dimensions.

Hsiang minimal cones

We shall concern with *cubic* minimal cones, i.e. cones of the smallest nontrivial degree 3.

Hsiang's problem (1967). How to characterize cubic minimal cones? Can one at least characterize all cubic polynomial solutions of

$$|Du|^2 \Delta u - \frac{1}{2} \langle Du, D|Du|^2 \rangle = \lambda |x|^2 u,$$

or, more explicitly,

$$(u_{x_1}^2 + \dots + u_{x_n}^2)(u_{x_1 x_1} + \dots + u_{x_n x_n}) - \sum_{i,j=1}^n u_{x_i} u_{x_j} u_{x_i x_j} = \lambda(x_1^2 + \dots + x_n^2) \cdot u(x)$$

Some explicit examples

- $u = \operatorname{Re}(z_1 z_2) z_3$, the triality polynomials in \mathbb{R}^{3d} , where $z_i \in \mathbb{R}, \mathbb{C}, \mathbb{H}$ or \mathbb{O}

- $$u(x) = \begin{vmatrix} \frac{1}{\sqrt{3}}x_1 + x_2 & x_3 & x_4 \\ x_2 & \frac{-2}{\sqrt{3}}x_1 & x_5 \\ x_4 & x_5 & \frac{1}{\sqrt{3}}x_1 - x_2 \end{vmatrix}$$

the generic norm in the Jordan algebra of 3×3 symmetric matrices over \mathbb{R}

- $$u(x) = \begin{vmatrix} x_1 & x_2 & x_3 \\ x_4 & x_5 & x_6 \\ x_7 & x_8 & x_9 \end{vmatrix}$$

the generic norm in the Jordan algebra of 4×4 symmetric *traceless* matrices over \mathbb{R}

Thus, Hsiang cubics are nicely encoded by certain algebraic structures. Which ones?

A metrized algebra approach

Let u be a cubic form on a vector space V over \mathbb{F} with an inner product $\langle x, y \rangle$. Linearize u to get a **symmetric trilinear** form $\tilde{u}(x, y, z)$ such that

$$\tilde{u}(x, x, x) = 6u(x)$$

Define a multiplication $(x, y) \rightarrow xy$ on V by the duality:

$$\tilde{u}(x, y, z) = \langle xy, z \rangle, \quad \forall z \in V.$$

Thus obtained algebra $V(u)$ is **commutative** and metrized:

$$\langle xy, z \rangle = \langle x, yz \rangle$$

and the cubic form is recovered by $u(x) = \frac{1}{6}\langle x, x^2 \rangle$. We also have

- $x^2 = 2 \operatorname{grad} u(x)$
- x is an idempotent in V iff kx is a stationary point of u ($x \parallel \operatorname{grad} u(x)$).
- $xy = \operatorname{Hess} u(x)y \Rightarrow \operatorname{ad}_x = \operatorname{Hess} u(x)$
- ad_x is **self-adjoint**: $\langle \operatorname{ad}_x y, z \rangle = \langle y, \operatorname{ad}_x z \rangle$
- Any idempotent is semisimple!

Algebras of minimal cones

In the above setup,

$$\langle Du, Du \rangle \Delta u - \langle Du, \text{Hess}(u) Du \rangle = \lambda |x|^2 u$$

yields

$$\left\langle \frac{x^2}{2}, \frac{x^2}{2} \right\rangle \text{tr}(\text{ad}_x) - \frac{1}{2} \left\langle \frac{x^2}{2}, \frac{x^3}{2} \right\rangle = \frac{\lambda}{6} \langle x, x \rangle \langle x, x^2 \rangle.$$

One can show that u is harmonic (non-trivial!), i.e.

$$\text{ad}_x \quad \text{is trace less!}$$

This implies that to any cubic cone one can attach a commutative metrized algebra with

$$\begin{cases} \langle x^2, x^3 \rangle &= \langle x, x \rangle \langle x^2, x \rangle \\ \text{tr}(\text{ad}_x) &= 0. \end{cases} \quad (1)$$

Conversely: if A satisfies (1) then $u(x) = \frac{1}{6} \langle x, x^2 \rangle$ generates a minimal cone.

We arrive at classification of all algebras satisfying (1). We call these **Hsiang algebras**.

Two basic examples in dimensions 2 and 3

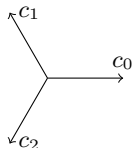
Example 1. Let A be the algebra on \mathbb{R}^2 generated by **idempotents** c_0, c_1, c_2 with

$$c_0 + c_1 + c_2 = 0$$

Then for any triple $\{i, j, k\} = \{1, 2, 3\}$

$$c_k = c_k^2 = (-c_i - c_j)^2 = c_i + c_j + 2c_i c_j = -c_k + 2c_i c_j$$

$$\Rightarrow c_i c_j = c_k \quad \Rightarrow c_k(c_i - c_j) = -1(c_i - c_j).$$



This implies that A is a Hsiang algebra, $A = A_1 \oplus A_{-1}$ with \mathbb{Z}_2 -graded fusion laws

\star	1	-1
1	1	-1
-1	-1	1

The **minimal cone** is given by $x_1^2 x_2 = 0$, i.e. pair of two orthogonal lines in \mathbb{R}^2

Two basic examples in dimensions 2 and 3

Example 2. Let A be generated by **idempotents** c_0, c_1, c_2, c_3 in \mathbb{R}^3 subject to

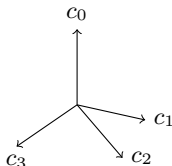
$$(c_i + c_j)^2 = 0 \quad \text{for } i \neq j$$

Then similarly to the above, one easily verifies that

$$A = A_1 \oplus A_{-\frac{1}{2}}, \quad \dim A_1 = 1, \quad \dim A_{-\frac{1}{2}} = 2.$$

Then A is a Hsiang algebra with fusion laws

\star	1	$-\frac{1}{2}$
1	1	$-\frac{1}{2}$
$-\frac{1}{2}$	$-\frac{1}{2}$	1, $-\frac{1}{2}$



After a *rank one perturbation* A becomes a Jordan algebra of Clifford type.

The **minimal cone** is given by $x_1 x_2 x_3 = 0$, i.e. the triple of coordinate planes in \mathbb{R}^3 .

Generalization

Definition. A commutative algebra V with associating form $\langle x, y \rangle$ is called **polar** if

- there is a \mathbb{Z}_2 -grading $V = V_0 \oplus V_1$,
- $V_0 V_0 = 0$,
- $x(xy) = |x|^2 y$ for $x \in V_0$ and $y \in V_1$.

Any polar algebra is Hsiang!

Remark. Can be simply constructed using **symmetric Clifford systems**, i.e. symmetric matrices

$$A_i^2 = I, \quad A_i A_j + A_j A_i = 0, \quad i \neq j$$

This yields an obstruction:

$$q \leq 1 + \rho(p),$$

where $\rho(m) = 8a + 2^b$, if $m = 2^{4a+b} \cdot \text{odd}$, $0 \leq b \leq 3$ is the **Hurwitz-Radon function**.

Definition. A Hsiang algebra V isomorphic to a polar algebra is of **Clifford type**; otherwise it is **exceptional**.

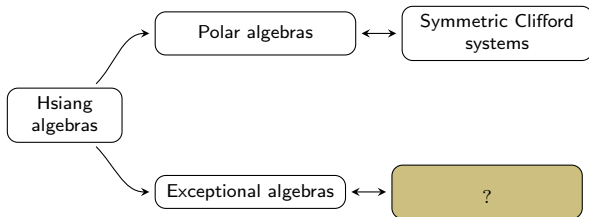
Exceptional Hsiang algebras

Any commutative **pseudocomposition algebra**, i.e. an algebra with

$$x^3 = |x|^2 x,$$

$$\text{tr ad}_x = 0$$

is an *exceptional* Hsiang algebra.



We prove that there are only finitely many exceptional Hsiang algebras!

Regular vs Exceptional

- Many classification results reminiscence the famous Cartan-Killing classification: there are several 'nice' series and finitely many exceptional (sporadic) elements.
 - ▶ the A , B , C , D -series of simple Lie groups and the exceptional Lie groups G_2 , F_4 , E_6 , E_7 , and E_8 ;
 - ▶ special Jordan algebras (obtained from associative algebras) vs the Albert 27-dimensional exceptional Jordan algebra
 - ▶ the classification of finite simple groups
- The fundamental question is why exceptional (sporadic) objects do really exist?

Finiteness of exceptional Hsiang algebras

The Peirce decomposition Let c be an idempotent. Then

$$A = A_1 \oplus A_{-1} \oplus A_{-\frac{1}{2}} \oplus A_{\frac{1}{2}}$$

where $A_1 = \mathbb{R}c$ (i.e. c is primitive!) satisfying the Hsiang algebras fusion laws:

\star	1	-1	- $\frac{1}{2}$	$\frac{1}{2}$
1	1	-1	- $\frac{1}{2}$	- $\frac{1}{2}$
-1	-1	1	$\frac{1}{2}$	- $\frac{1}{2}$, $\frac{1}{2}$
- $\frac{1}{2}$	- $\frac{1}{2}$	$\frac{1}{2}$	1, - $\frac{1}{2}$	-1, $\frac{1}{2}$
$\frac{1}{2}$	$\frac{1}{2}$	- $\frac{1}{2}$, $\frac{1}{2}$	-1, $\frac{1}{2}$	1, -1, - $\frac{1}{2}$

- Two distinguished subalgebras: $A_1 \oplus A_{-1}$ (carries a hidden Clifford algebra structure) and $A_1 \oplus A_{-\frac{1}{2}}$ (carries a hidden rank 3 Jordan algebra structure)
- Primitive idempotents w in the hidden Jordan algebra ($w \bullet w = w$) are exactly 2-nilpotents in V ($w^2 = w$) with the fusion rules

Finiteness of exceptional Hsiang algebras

Main Theorem (V.T., 2014, 2016)

(i) Given an idempotent $c \in A$, the new algebra on $\Lambda_c = (A_1 \oplus A_{-\frac{1}{2}}, \bullet)$ with

$$x \bullet y = \frac{1}{2}xy + \langle x, c \rangle y + \langle y, c \rangle x - 2\langle xy, c \rangle c.$$

is a Euclidean Jordan algebra of rank 3 with unit $c^* = 2c$.

(ii) A is exceptional if and only if the Jordan algebra Λ_c is simple.

(iii) There holds

$$\dim A_{-1} - 1 \leq \rho(\dim A_{-1} + \dim A_{-\frac{1}{2}} - 1),$$

where ρ is the Hurwitz-Radon function.

(iv) There are at most 16 possible classes of exceptional Hsiang algebras:

$\dim A$	2	5	8	14	26	9	12	21	15	18	24	30	42	27	30	54
$\dim A_{-1}$	1	2	3	5	9	0	1	4	0	1	3	5	9	0	1	1
$\dim A_{-1/2}$	0	0	0	0	0	5	5	5	8	8	8	8	8	14	14	26

unsettled

Realization of cubic cones via Jordan algebras

Theorem (Existence)

Denote by $V(u)$ the algebra on a vector inner product space V generated by the cubic form u . Then $A = V(u)$ where

- $\dim A_{-\frac{1}{2}} = 0$, $\dim A_{-1} = d + 1$, and $u = \frac{1}{6}\langle z, z^2 \rangle$, $V = \mathcal{H}_3(\mathbb{K}_d) \ominus \mathbb{R}e$,
 $d = 0, 1, 2, 4, 8$.
- $\dim A_{-1} = 0$, $\dim A_{-\frac{1}{2}} = \frac{3}{2}d + 2$, and $u = \frac{1}{12}\langle z^2, 3\bar{z} - z \rangle$, where $z \rightarrow \bar{z}$ is the natural involution on $V = \mathcal{H}_3(\mathbb{K}_d)$, $d = 2, 4, 8$.
- $\dim A_{-1} = 1$, $\dim A_{-\frac{1}{2}} = 3d + 2$, and $u(z) = \operatorname{Re}\langle z, z^2 \rangle$, where $z \in V = \mathcal{H}_3(\mathbb{K}_d) \otimes \mathbb{C}$, $d = 1, 2, 4, 8$.
- $\dim A_{-1} = 4$ and $\dim A_{-\frac{1}{2}} = 5$, where $u = \frac{1}{6}\langle z, z^2 \rangle$ on $\mathcal{H}_3(\mathbb{K}_8) \ominus \mathcal{H}_3(\mathbb{K}_1)$

Epilogue: Nonassociative algebras and singular solutions

- Evans, Crandall, Lions, Jensen, Ishii: If $\Omega \subset \mathbb{R}^n$ is bounded with C^1 -boundary, ϕ continuous on $\partial\Omega$, F uniformly elliptic operator then the Dirichlet problem

$$\begin{aligned}F(D^2u) &= 0, \quad \text{in } \Omega \\ u &= \phi \quad \text{on } \partial\Omega\end{aligned}$$

has a unique **viscosity solution** $u \in C(\Omega)$;

- Krylov, Safonov, Trudinger, Caffarelli, early 80's: the solution is always $C^{1,\varepsilon}$
- Nirenberg, 50's: if $n = 2$ then u is classical (C^2) solution
- Nadirashvili, Vlăduț, 2007-2011: if $n \geq 12$ then there are solutions which are not C^2 .

Theorem (Nadirashvili-V.T.-Vlăduț, 2012, 2015) *The function $w(x) := \frac{u_1(x)}{|x|}$ where*

$$u_1(x) = x_5^3 + \frac{3}{2}x_5(x_1^2 + x_2^2 - 2x_3^2 - 2x_4^2) + \frac{3\sqrt{3}}{2}x_4(x_2^2 - x_1^2) + 3\sqrt{3}x_1x_2x_3,$$

is a singular viscosity solution of the uniformly elliptic Hessian equation

$$(\Delta w)^5 + 2^8 3^2 (\Delta w)^3 + 2^{12} 3^5 \Delta w + 2^{15} \det D^2(w) = 0,$$

This also gives the best possible dimension ($n = 5$) where homogeneous order 2 real analytic functions in $\mathbb{R}^n \setminus \{0\}$.

Some important questions remain unanswered

How incident (important, relevant) that the certain *commutative non-associative* algebraic structures coming from finite simple groups, geometry of minimal cones and PDEs (truly viscosity solutions)

- have a distinguished Peirce spectrum
- have distinguished (in particular, graded) fusion rules
- are axial
- are metrized (i.e. carrying an associating symmetric bilinear form)?

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THANK YOU FOR YOUR ATTENTION!